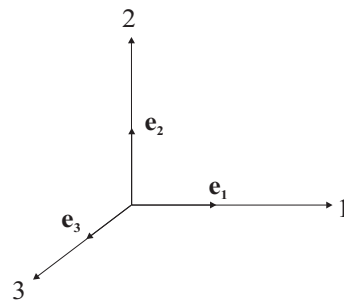


Handout # 3: Cartesian Tensors

Vectors in Cartesian Coordinates

Coordinates system E with unit vectors

$$\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3 \Rightarrow \mathbf{e}_i \text{ (vector in } i\text{-direction)}$$



Since unit vectors are orthogonal basis

$$\mathbf{e}_i \cdot \mathbf{e}_j = \delta_{ij}$$

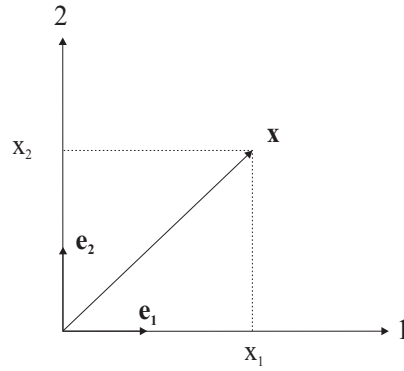
with Kronecker delta δ_{ij} defined by

$$\begin{aligned} \delta_{ij} &= \begin{cases} 1 & \text{if } i = j \\ 0 & \text{if } i \neq j \end{cases} \\ &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \end{aligned}$$

Index substitution:

$$\delta_{ij} x_i = x_j$$

Coordinate System in 2D



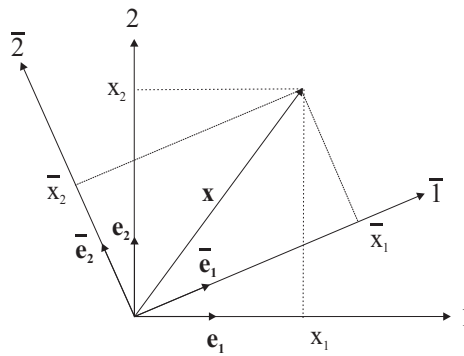
Any vector \boldsymbol{x} can be represented in E by

$$\begin{aligned}\boldsymbol{x} &= \boldsymbol{e}_1 x_1 + \boldsymbol{e}_2 x_2 + \boldsymbol{e}_3 x_3 \\ &= \boldsymbol{e}_i x_i \quad (\text{dummy index implies summation})\end{aligned}$$

Definition of Tensor

- **Tensor consists of tensor components (e.g. x_i) and coordinate system (e.g. E)**
- Tensor remains unchanged if expressed in coordinate system \bar{E} obtained by axes rotation and reflection of system E

Consider two coordinate systems E and \bar{E} :



$$\mathbf{x} = \mathbf{e}_i x_i = \overline{\mathbf{e}_i} \overline{x_i}$$

$$\mathbf{e}_k \cdot \mathbf{e}_i x_i = \overline{\mathbf{e}_k} \cdot \overline{\mathbf{e}_i} \overline{x_i}$$

$$\delta_{ki} x_i = x_k = \overline{a_{ki}} \overline{x_i}$$

where

$$a_{ki} = \mathbf{e}_k \cdot \overline{\mathbf{e}_i} = |\mathbf{e}_k| |\overline{\mathbf{e}_i}| \cos \alpha_{ki} = \cos \alpha_{ki}$$

since $|\mathbf{e}_k| = |\overline{\mathbf{e}_i}| \equiv 1$.

If transformation can be written as

$$x_i = \overline{a_{ij}} \overline{x_j}$$

then \mathbf{x} is a tensor.

Second order tensor

$$\mathbf{b} = \mathbf{e}_i \mathbf{e}_j b_{ij} = \overline{\mathbf{e}_i} \overline{\mathbf{e}_j} \overline{b_{ij}}$$

Tensor operations

- Tensor product

$$\mathbf{u} \mathbf{b} = \mathbf{e}_i u_i \mathbf{e}_j \mathbf{e}_k b_{jk} = \mathbf{e}_i \mathbf{e}_j \mathbf{e}_k u_i b_{jk} \quad (\text{order } N + M)$$

- Inner product

$$\mathbf{u} \cdot \mathbf{b} = \mathbf{e}_i u_i \cdot \mathbf{e}_j \mathbf{e}_k b_{jk} = \delta_{ij} \mathbf{e}_k u_i b_{jk} = \mathbf{e}_k u_j b_{jk} \quad (\text{order } N + M - 2)$$

- Gradient

$$\nabla \equiv \mathbf{e}_i \frac{\partial}{\partial x_i}$$

$$\nabla \mathbf{b} = \mathbf{e}_i \frac{\partial}{\partial x_i} (\mathbf{e}_j \mathbf{e}_k b_{jk}) = \mathbf{e}_i \mathbf{e}_j \mathbf{e}_k \frac{\partial b_{jk}}{\partial x_i} \quad (\text{order } 1 + M)$$

- Divergence

$$\nabla \cdot \mathbf{b} = \mathbf{e}_i \cdot \frac{\partial}{\partial x_i} (\mathbf{e}_j \mathbf{e}_k b_{jk}) = \delta_{ij} \mathbf{e}_k \frac{\partial b_{jk}}{\partial x_i} = \mathbf{e}_k \frac{\partial b_{ik}}{\partial x_i} \quad (\text{order } M-1)$$

- Cross product

$$\mathbf{u} \times \mathbf{v} = \det \begin{vmatrix} \mathbf{e}_1 & \mathbf{e}_2 & \mathbf{e}_3 \\ u_1 & u_2 & u_3 \\ v_1 & v_2 & v_3 \end{vmatrix}$$

Alternating symbol:

$$\epsilon_{ijk} = \begin{cases} 1 & \text{if } ijk \text{ are cyclic} & (123, 231, 312) \\ -1 & \text{if } ijk \text{ are anti-cyclic} & (321, 213, 132) \\ 0 & \text{otherwise} \end{cases}$$

$$\Rightarrow \mathbf{u} \times \mathbf{v} = \epsilon_{ijk} \mathbf{e}_i u_j v_k \quad \text{and} \quad \boldsymbol{\omega} = \nabla \times \mathbf{u} = \epsilon_{ijk} \mathbf{e}_i \frac{\partial u_k}{\partial x_j}$$

Note: ϵ_{ijk} is not a tensor, hence the cross product is not a tensor operation.

Momentum equations in direct and suffix notation:

$$\frac{\partial \mathbf{U}}{\partial t} + \underbrace{\mathbf{U} \cdot \nabla \mathbf{U}}_{\nabla \cdot (\mathbf{U} \mathbf{U})} = -\frac{1}{\rho} \nabla P + \nu \underbrace{\nabla \cdot \nabla \mathbf{U}}_{\nabla^2 \mathbf{U}}$$

$$\frac{\partial U_j}{\partial t} + \underbrace{U_i \frac{\partial U_j}{\partial x_i}}_{\frac{\partial U_i U_j}{\partial x_i}} = -\frac{1}{\rho} \frac{\partial P}{\partial x_j} + \nu \frac{\partial^2 U_j}{\partial x_i^2}$$